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SUNSET OVER BROWNISTAN*

by

ERHAN ÇINLAR

Princeton University
Department of Civil Engineering and Operations Research
School of Engineering and Applied Science
Princeton, New Jersey 08544

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Abstract

Consider a Brownian motion with a downward drift of rate a . Its maximum over all time has the exponential distribution with parameter $2a$. Our aim is to study this maximum as a stochastic process indexed by a . That process is related to the convex majorant of the standard Brownian motion and, through the latter, to a Poisson random measure. This connection is exploited to obtain various distributional results. The results are of interest in queueing theory.

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considered by NEWELL [4] and by COFFMAN, KADOTA, and SHEPP [2], the latter viewing the model as that of storage allocation in computer memory.

Let $Q_i(n)$ be the random variable that is 0 or 1 according as the stall n is empty or occupied at time t . The random vector $Q_i = (Q_i(1), Q_i(2), \dots)$ is the state of the system at time t . The process $\{\sum_n Q_i(n); t \geq 0\}$ is the queue size process in an M/M/ ∞ system; it is regenerative, and 0 is a regeneration state for it. It follows that the vector $(0, 0, \dots)$ is a regeneration state for $\{Q_i; t \geq 0\}$ and that the latter has an equilibrium distribution. Let Q be a random vector (of zeros and ones) whose law is that equilibrium distribution.

The distribution of $\sum_n Q(n)$ is Poisson with mean λ . The distribution of $\sum_1^m Q(n)$ is the equilibrium distribution of the queue size process in the M/M/m/m system with arrival rate λ and service rate 1; thus, that distribution is the conditional distribution of $\sum_n Q(n)$ given that $\sum_n Q(n) \leq m$. Other than these facts and a few conclusions that can be drawn from them by elementary probabilistic considerations, there is not much known about the distribution of Q .

For $a > 0$, let $\lambda^{1/4} Y_\lambda(a, t)$ be the number of empty stalls at time t among those labeled with $n < \lambda - a\lambda^{3/4}$. ALDOUS [1] has shown that the process $\{Y_\lambda(a, t); a > 0, t \geq 0\}$ converges weakly, as $\lambda \rightarrow \infty$, to a process $\{Y(a, t); a > 0, t \geq 0\}$, which he identified and showed that, in the limit as $t \rightarrow \infty$, converges weakly to the process

$$Y(a) = \max_{t \geq 0} (\sqrt{2} B_t - at), \quad a > 0,$$

where B is the standard Brownian motion. He calls $\{Y(a); a > 0\}$ the exponential process, after the well-known fact that $Y(a)$ has the exponential distribution with mean $1/a$ for each a .

Our main contribution is to supply the probability law of Y in simpler terms. For this purpose we choose to work with

$$(1.1) \quad Z_a = \frac{1}{\sqrt{2}} Y(\sqrt{2} a) = \max_{t \geq 0} (B_t - at), \quad a > 0,$$

and let D_a be the last time t at which $B_t - at$ touches its zenith Z_a , that is,

$$(1.2) \quad D_a = \sup \{t : B_t - at = Z_a\}, \quad a > 0.$$

It turns out that D_a is the left-derivative of Z at a and, thus, is related to the density of empty stalls in the parking lot, in equilibrium, around $\lambda - a\lambda^{3/4}$ for large λ .

The next section contains a few simple geometric observations. First we relate the process (D, Z) to the convex majorant of the Brownian motion B . Using the hard results of GROENEBOOM [3] and PITMAN [5] about the latter, we are able to express D and Z in terms of a Poisson random measure on $(0, \infty) \times (0, \infty)$. It follows, in particular, that D has non-stationary independent increments, and that (D, Z) is a non-homogeneous Markov process.

The process $a \rightarrow Z_a$ is continuous, concave, and decreases from its limit $+\infty$ at $a = 0+$ to its limit 0 at ∞ . Therefore, its "hitting time" process

$$(1.3) \quad A_z = \inf \{a : Z_a < z\}, \quad z > 0,$$

is the functional inverse of Z . It turns out that the process A has the same probability law as Z . This observation is also put in the next section.

The last section is devoted to computational issues. We compute the joint distribution of D_a and Z_a and also the transition function of the Markov process (D, Z) .

2. ZENITH PROCESS

The problem with the definition (1.1) of Z_a is that it suggests re-drawing the path $t \rightarrow B_t - at$ if we wish to vary a . The following observation circumvents the problem:

$$(2.1) \quad Z_a = \inf \{x > 0 : x + at > B_t \text{ for all } t \geq 0\}.$$

Obviously, this is a re-wording of (1.1), but the mental picture it suggests is much more convenient for manipulating a : the line $t \rightarrow Z_a + at$ is the infimum of all lines of slope a that never touch B . This picture is drawn in Figure 1 below.

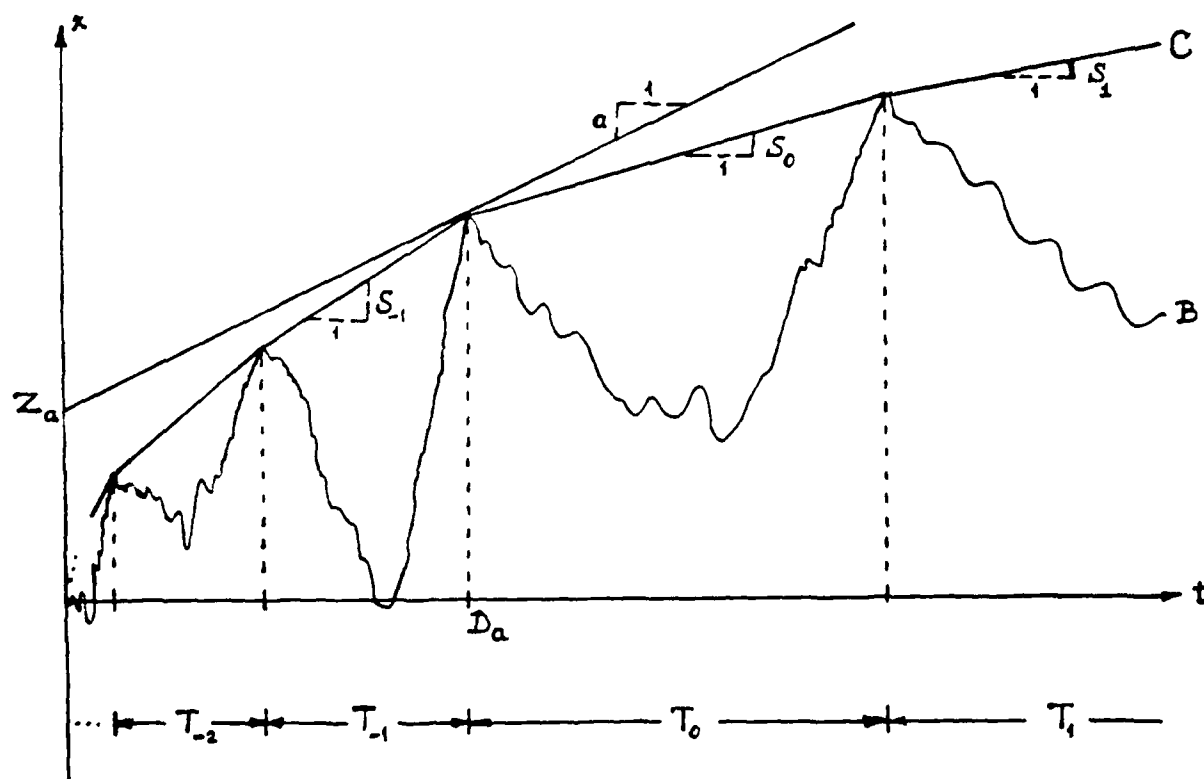


FIGURE 1

Let C denote the convex majorant of B , that is, the minimal convex path that dominates B (see Figure 1 again). Note that Z_a and D_a are determined by C : the line $t \rightarrow Z_a + at$ is the infimum of all lines of slope a that never touch C , and D_a is the last time at which that line touches C . In fact, for fixed $a > 0$, D_a is almost surely the only time t with $C_t = Z_a + at$.

It is known (see GROENEBOOM [3] for instance) that C is continuous and piecewise linear. The countable collection of its vertices has, almost surely, only one accumulation point, namely $(0,0)$. Fix an $a > 0$; note that $(D_a, Z_a + aD_a)$ is a vertex; let T_0, T_1, \dots be the lengths of successive intervals of linearity going to the right from D_a ; let T_{-1}, T_{-2}, \dots be those to the left; and let S_i be the slope of C over the interval whose length is denoted by T_i . The following major result was obtained by GROENEBOOM [3]; a simpler proof using the excursions of B may be found in PITMAN [5].

(2.2) THEOREM. The pairs (S_i, T_i) , $i \in \mathbb{Z}$, form a Poisson random measure N on $(0, \infty) \times (0, \infty)$ whose mean measure is

$$(2.3) \quad \nu(ds, dt) = \frac{ds}{s} \gamma_s(dt),$$

where γ_s is the gamma distribution with shape index $1/2$ and scale parameter $s^2/2$ (the corresponding mean is $1/s^2$).

The probability law of a Poisson random measure is determined by its mean measure. Thus, the following specifies the probability law of (D, Z) . For computational purposes, the representations given here for D_a and Z_a are the key starting points.

(2.4) PROPOSITION. For each $a > 0$,

$$(2.5) \quad D_a = \int_{[a, \infty) \times (0, \infty)} N(ds, dt) t,$$

$$(2.6) \quad Z_a = \int_a^\infty D_s ds = \int_{[a, \infty) \times (0, \infty)} N(ds, dt) (s - a) t.$$

The process D has non-stationary independent increments. The process (D, Z) is a temporally non-homogeneous Markov process.

PROOF. First note that (see Figure 1)

$$D_a = \sum_i T_i 1_{[a, \infty)}(S_i),$$

$$Z_a + a D_a = B(D_a) = \sum_i S_i T_i 1_{[a, \infty)}(S_i).$$

Expressed in terms of the Poisson random measure N , these become (2.5) and (2.6). The remaining statements are immediate from the independence of the restrictions of N to disjoint Borel sets.

Figure 2 below shows the qualitative features of D : it is piecewise constant, left-continuous, and decreases from its limit $+\infty$ at $a = 0+$ to its limit 0 at $+\infty$. It follows from (2.6) that Z is continuous, concave, piecewise linear, and decreases from its limit $+\infty$ at $a = 0+$ to 0 at $+\infty$.

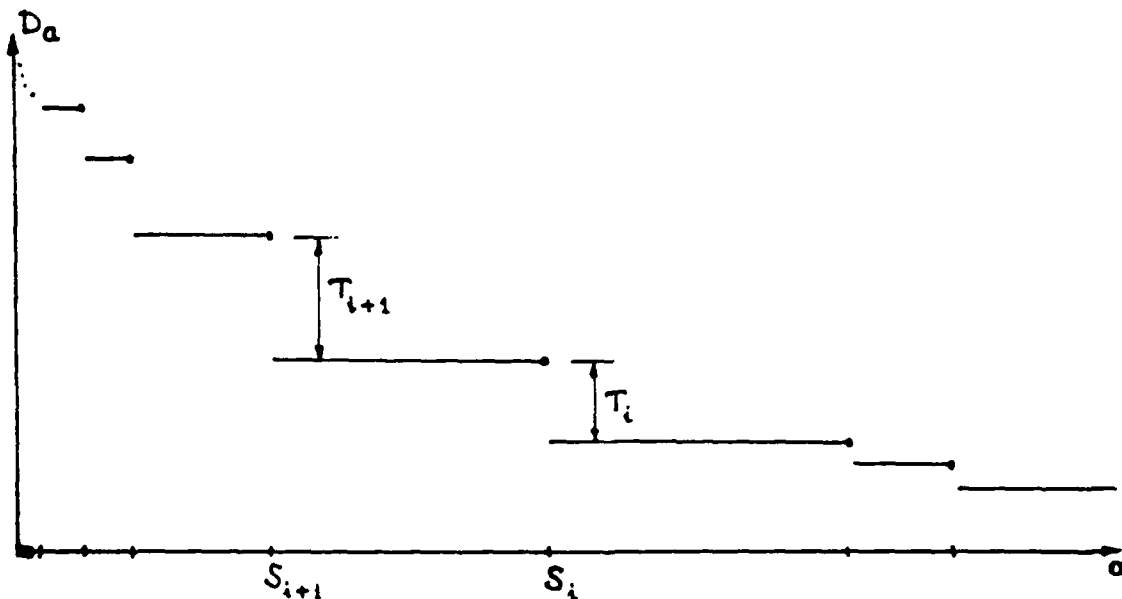


FIGURE 2

It was noted by ALDOUS [1] that, for each $c > 0$, the process $(c^2 D_{ca}, cZ_{ca})_{a>c}$ has the same probability law as (D, Z) . This can be seen from the preceding characterization: $a \rightarrow c^2 D_{ca}$ jumps at the points S_i/c by the amounts $c^2 T_i$; the pairs $(S_i/c, c^2 T_i)$ form a Poisson random measure that has the same mean measure as N ; hence, $a \rightarrow c^2 D_{ca}$ has the same law as D .

We end this section with an observation on the process

$$(2.7) \quad A_z = \inf\{a : Z_a < z\}, \quad z > 0.$$

Obviously, $z \rightarrow A_z$ is the functional inverse of the one-to-one mapping $a \rightarrow Z_a$ of $(0, \infty)$ onto $(0, \infty)$. It follows that the qualitative picture of A is exactly that of Z . In particular, A is piecewise linear and

$$(2.8) \quad \hat{D}_z = \lim_{\epsilon \downarrow 0} \frac{A_{z+\epsilon} - A_z}{\epsilon} = \frac{1}{D(A_z)}, \quad z > 0.$$

The process \hat{D} is piecewise constant, right-continuous, decreasing.

(2.9) PROPOSITION. The process (\hat{D}, A) has the same probability law as the process (D, Z) . In particular, the collection $\{Z(S_i); i \in \mathbb{Z}\}$ has the same law as the collection $\{S_i; i \in \mathbb{Z}\}$; they form Poisson random measures on $(0, \infty)$ with mean measure ds/s .

PROOF. We put (2.7) and (2.1) together and manipulate:

$$\begin{aligned} A_z &= \inf\{a : \inf\{x : x + at > B_t \text{ for all } t\} < z\} \\ &= \inf\{a : z + at > B_t \text{ for all } t\} \\ &= \inf\{a : a + zu > u B_{1/u} \text{ for all } u\}. \end{aligned}$$

This shows that A is the zenith process associated with the process $\{uB_{1/u}; u \geq 0\}$, just as Z is

the zenith process associated with B . Since $(\mu B_{1/\mu})$ is a standard Brownian motion like B , it follows that A has the same probability law as Z . This proves the first statement, since \hat{D} is the derivative of $-A$ and D is the derivative of $-Z$.

The points S_i are the jump locations of D , and the points $Z(S_i)$ are those of \hat{D} . This proves the second statement.

3. ENTRANCE LAW AND TRANSITION FUNCTION

We derive the distribution of the random variable (D_a, Z_a) and the transition function of the process (D, Z) . The computations rest on the characterization given by Proposition (2.4) and on the well-known formula for the Laplace functional of the Poisson random measure N with mean measure ν :

$$(3.1) \quad E \exp - \int N(dx) f(x) = \exp - \int \nu(dx) (1 - e^{-f(x)})$$

for every positive Borel function f on $(0, \infty) \times (0, \infty)$.

(3.2) PROPOSITION. For each $a > 0$,

$$(3.3) \quad E \exp (-pD_a - qZ_a) = \frac{2a}{a + q + \sqrt{a^2 + 2p}}, \quad p \geq 0, q \geq 0;$$

$$(3.4) \quad P\{D_a \in dt, Z_a \in dz\} = dt dz \frac{2az}{\sqrt{2\pi t^3}} \frac{e^{-(z+at)^2/2t}}{\sqrt{2\pi t^3}}, \quad t > 0, z > 0.$$

In particular,

$$(3.5) \quad P\{Z_a \in dz\} = dz 2ae^{-2az}, \quad P\{D_a \in dt\} = dt \int_0^\infty du \frac{ae^{-a^2u/2}}{\sqrt{2\pi u^3}}.$$

PROOF. Fix $a > 0, p \geq 0, q \geq 0$. In view of (2.5) and (2.6),

$$pD_a + qZ_a = \int_{[a, \infty) \times (0, \infty)} N(ds, dt) (pt + q(s-a)t).$$

Using (3.1) and the form of the mean measure ν given by (2.3), the Laplace transform (3.3) is obtained via elementary calculus. To invert the Laplace transform, first write it as

$$\int_0^\infty dz e^{-qz} 2ae^{-az} e^{-z\sqrt{a^2+2p}}$$

and then recall that $e^{-z\sqrt{2r}}$, $r \geq 0$, is the Laplace transform of H_z , the first time a standard Brownian motion hits the level z , that is,

$$e^{-z\sqrt{a^2+2p}} = \int_0^\infty dt \cdot \frac{z e^{-z^2/2t}}{\sqrt{2\pi t^3}} e^{-(p+a^2/2)t}.$$

The rest is trivial.

(3.6) REMARK. Although (3.4) is explicit and shades of exponential and stable distributions can be felt, it does not seem well-suited for probabilistic thinking. The following representation is better, especially for Monte-Carlo methods. For $a > 0$,

$$a^2 D_a = X (1 - \sqrt{U})^2, \quad aZ_a = X \sqrt{U} (1 - \sqrt{U}),$$

where X and U are independent, U has the uniform distribution on $(0,1)$, and X has the gamma distribution with shape index $3/2$ and scale parameter $1/2$.

The following specifies the joint Laplace transform of any number of increments of Z (upon taking $f = p_1 1_{A_1} + \dots + p_n 1_{A_n}$ with A_1, \dots, A_n disjoint intervals).

(3.7) PROPOSITION. For any positive Borel function f on $(0, \infty)$,

$$E \exp \int_{(0, \infty)} f(a) dZ_a = \exp - \int_0^\infty ds \left(\frac{1}{s} - \frac{1}{\sqrt{s^2 + 2\bar{f}(s)}} \right)$$

where $\bar{f}(s)$ is the Lebesgue integral of f over $(0, s)$.

PROOF. Note that

$$\int f(a) dZ_a = - \int f(a) D_a da = - \int N(ds, dt) \bar{f}(s)t ,$$

and use (3.1).

As mentioned in Proposition (2.4), the process D has non-stationary independent increments, and the process (D, Z) is a non-homogeneous parameter Markov process. Let

$$(3.8) \quad P_{ab}(t, x; du, dy) = P\{D_b \in du, Z_b \in dy \mid D_a = t, Z_a = x\}$$

for $0 < b < a$, $0 < t < u$, $0 < x < y$ (in our zeal to deal with positive random variables, we choose to work with the parameters in decreasing order). In view of (2.5) and (2.6),

$$(3.9) \quad P_{ab}(t, x; du, dy) = P\{t + U \in du, x + (a - b)t + Y \in dy\} ,$$

where

$$(3.10) \quad U = \int_{(b, a) \times (0, \infty)} N(ds, dt)t, \quad Y = \int_{(b, a) \times (0, \infty)} N(ds, dt)(s - b)t .$$

The joint Laplace transform of U and Y can be obtained from (3.1) as in the first step of the proof of (3.2):

$$\begin{aligned} (3.11) \quad E e^{-pU - qY} &= \frac{b}{a} \cdot \frac{a + q + \sqrt{a^2 + 2p + 2(a - b)q}}{b + q + \sqrt{b^2 + 2p}} \\ &= \frac{b}{a} + \frac{a - b}{a} \cdot \frac{1}{2} \cdot \frac{2b}{b + q + \sqrt{b^2 + 2p}} \\ &\quad + \frac{a - b}{a} \cdot \frac{1}{2} \cdot \frac{2b}{b + q + \sqrt{b^2 + 2p}} \cdot \frac{1}{a - b} \int_b^a \frac{(c + q) dc}{\sqrt{(c + q)^2 + 2p - 2bq - q^2}} . \end{aligned}$$

Inverting this is tedious but manageable. It yields the following for the distribution φ of the pair

(U, Y) :

$$(3.12) \quad \epsilon \varphi = \frac{b}{a} \delta_{(0,0)} + (1 - \frac{b}{a}) \left(\frac{1}{2} \lambda_b + \frac{1}{2} \lambda_b * \frac{\mu_{bb} - \mu_{ab}}{a - b} \right)$$

where the asterisque denotes convolution, δ_x is the Dirac measure at x , λ_b is the distribution of (D_b, Z_b) specified by (3.4), and

$$(3.13) \quad \mu_{ab}(dt, dy) = \frac{e^{-b^2 t/2} dt}{\sqrt{2\pi t^3}} \delta_{(a-b)t}(dy), \quad t > 0, y > 0.$$

Putting the distribution $\epsilon \varphi$ of (U, Y) into (3.9) yields an explicit expression for the transition function P_{ab} . As a by-product, we have the joint distribution of

$$U = D_b - D_a, \quad Y = Z_b - Z_a - (b - a) D_a.$$

Noting that D_a is independent of (U, Y) , one can obtain the distribution of $(D_b - D_a, Z_b - Z_a)$ among other things.

However, it is clear that such results are of limited use because of their complexity. Overall, the computational complexity is caused by a confluence of two incompatible operations, addition and multiplication: look at the form (2.3) of the mean measure ν ; the Haar measure ds/s indicates that the natural group operation on the jump points S_i is multiplication, whereas the jump amounts T_i are additive.

Of course, it is easy to transform the Poisson random measure N into one with a nicer intensity: define f to be the mapping $(s, t) \rightarrow (\log s, s^2 t)$; then the image of N under f is the Poisson random measure $\hat{N} = Nf^{-1}$ on $(-\infty, \infty) \times (0, \infty)$ with mean measure $du \gamma(dv)$ where γ is the gamma distribution with shape and scale parameters equal to $1/2$. But, then, expressions for D_a and Z_a in terms of \hat{N} have to undo the transformation, and there is no gain at the end. Using $p = \log a$ to index the processes involved (and working with $\hat{Z}_p = Z(e^p)$) does not help either.

On the other hand, it is easy to describe the construction of the path of (D, Z) over the interval $(0, a]$. This may be useful for Monte-Carlo purposes.

First, we observe that the conditional distribution of S_{i+1} given (S_i, S_{i-1}, \dots) is the uniform distribution on $(0, S_i)$. Thus, to construct (D, Z) over $(0, a]$, we start with U and X described in Remark (3.6) and generate D_a and Z_a . Then, we let U_1, U_2, \dots be i.i.d. uniform on $(0, 1)$, let X_1, X_2, \dots be i.i.d. Gaussian with mean 0 and variance 1, and set

$$(3.14) \quad S_0 = a, \quad S_i = a U_1 U_2 \cdots U_i, \quad T_i = \left(\frac{X_i}{S_i}\right)^2, \quad i = 1, 2, \dots$$

With these, define

$$(3.15) \quad D(S_0) = D_a, \quad D(S_i) = D(S_{i-1}) + T_i, \quad i \geq 1,$$

$$(3.16) \quad Z(S_0) = Z_a, \quad Z(S_i) = Z(S_{i-1}) + (S_{i-1} - S_i) \cdot D(S_{i-1}), \quad i \geq 1.$$

Then, $(D_b)_{b \in (0, a]}$ is the left-continuous piecewise constant path whose value at S_i is $D(S_i)$, and $(Z_b)_{b \in (0, a]}$ is the continuous piecewise linear path whose value at S_i is $Z(S_i)$. Incidentally, (S_i) , $(S_i, D(S_i))$, $(S_i, D(S_i), Z(S_i))$ are all Markov chains.

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